

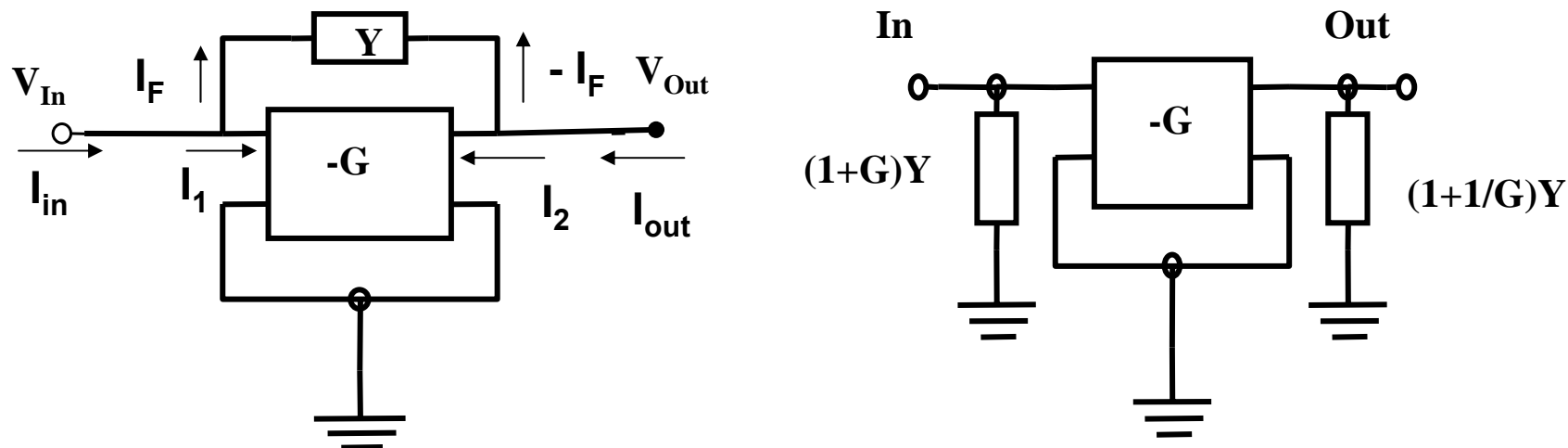
The Miller Theorem

and the

Frequency Response of the Common Emitter Amplifier

A useful approximation: The Miller Theorem

- Consider a shunt admittance connected between the input and output of an inverting voltage amplifier of gain G .



- Assuming that Y does not affect the gain, KCL at the input says that:

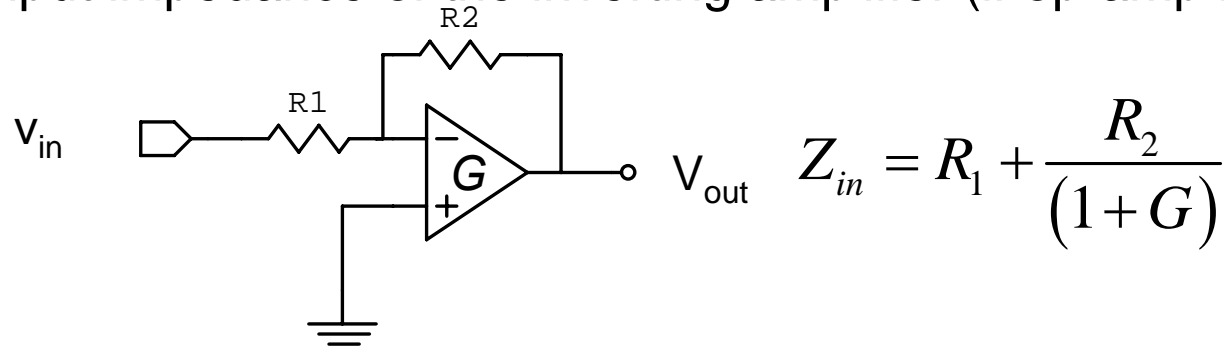
$$i_{in} = i_1 + i_F = i_1 + (v_{in} - v_{out})Y = i_1 + [(1+G)Y]v_{in}$$

$$i_{out} = i_2 - i_F = i_2 - (v_{in} - v_{out})Y = i_2 + [(1+1/G)Y]v_{out}$$

- This is the same current as we would get if we connected $Y(1+G)$ in parallel to the input and $Y(1+1/G)$ in parallel to the output (on the right)
- Note that this is an approximation; Y always changes the gain, unless $Z_{out} // Z_L = 0$. Nonetheless, very often it is a good approximation.**

An application of the Miller theorem

Input impedance of the inverting amplifier (if op-amp has zero Y_{in}):



We can even use the Miller theorem to calculate the gain of this circuit for finite op-amp gain G :

$$A_v = \frac{v_{out}}{v_{in}} = -G \frac{R_2 / (G+1)}{R_1 + R_2 / (G+1)} = \frac{-GR_2}{R_1G + R_1 + R_2}, \quad \lim_{G \rightarrow \infty} A_v = \frac{-R_2}{R_1}$$

This is the correct answer for the closed loop gain as we shall later see!

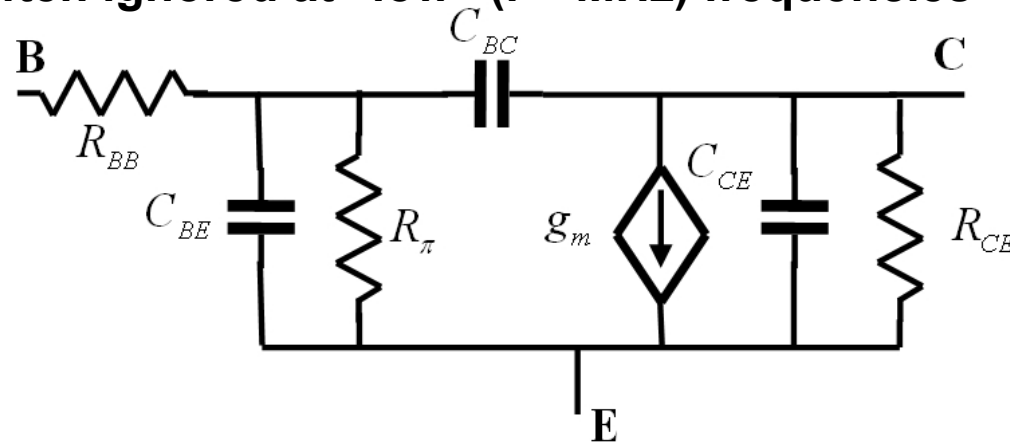
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The small signal model the of BJT

- The model below is useful to frequencies of up to 100s of MHz
- The model is strictly valid for B-E voltage variations smaller than $V_T=25\text{mV}$.
- To analyse a circuit:
 - *Each transistor in a circuit is replaced by this model*
 - *The “base spreading resistance” R_{BB} is lumped with the thevenin impedance of the signal source. It is not shown explicitly in these notes.*

What does R_{BB} do:

- R_{BB} is responsible for most of the noise of transistor amplifiers
- R_{BB} is responsible for the power gain rolling-off at high frequencies
- R_{BB} is often ignored at “low” ($f < \text{MHz}$) frequencies



BJT model: impedances and transconductance

- The large signal BJT equation gives the transconductance:

$$I_C = I_0 \left(e^{V_{be}/V_{th}} - 1 \right) \left(1 + \frac{V_{CE}}{V_A} \right) \approx I_0 e^{V_{be}/V_{th}}$$

$$V_{th} = \frac{k_B T}{e} \approx 26 \text{ mV at room temperature}$$

$$g_m = \frac{\partial I_C}{\partial V_{be}} \approx \frac{I_C}{V_{th}}$$

- However, since $I_E = (1 + 1/\beta) I_C$,

$$R_\pi = \frac{\partial V_{BE}}{\partial I_B} = \beta V_T / I_C = \beta / g_m$$

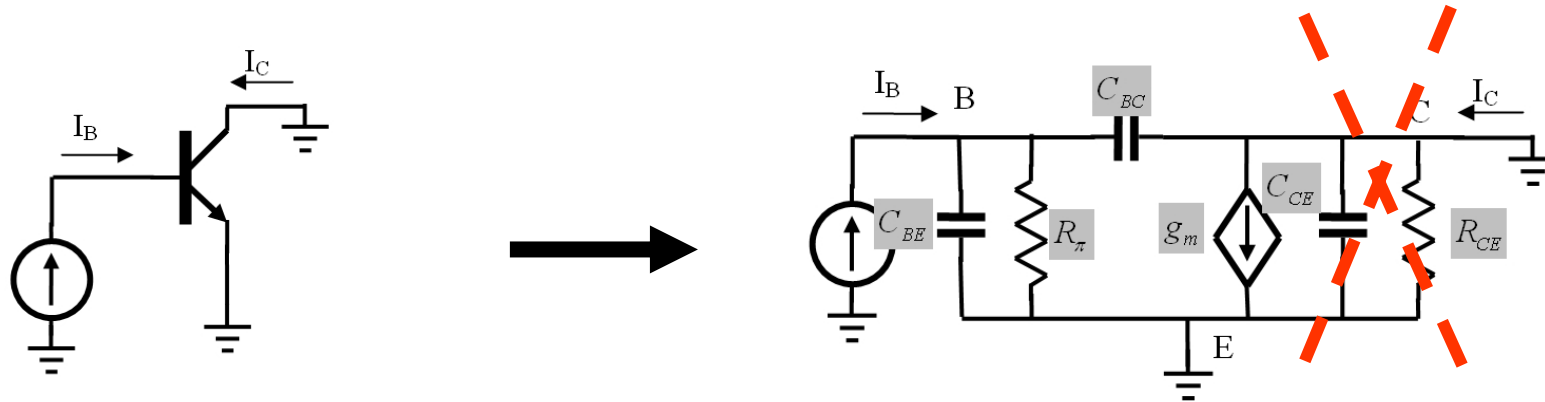
- Same equation gives output conductance:

$$G_{CE} = 1 / R_{CE} = \frac{\partial I_C}{\partial V_{CE}} = \frac{I_C}{V_A}$$

BJT model Capacitances

<ul style="list-style-type: none"> • BC junction is reverse biased: <ul style="list-style-type: none"> – <u>depletion capacitance</u> – <i>overlap capacitance</i> 	$C_{BC} = C_{BC0} + C_j$
<ul style="list-style-type: none"> • BE junction is forward biased: <ul style="list-style-type: none"> – <i>depletion capacitance</i> – <u>storage (diffusion capacitance)</u> – <i>overlap capacitance</i> 	$C_{BE} = C_{BE0} + C_j + C_D$
<ul style="list-style-type: none"> • CE capacitance is small <ul style="list-style-type: none"> – <i>overlap capacitance</i> 	$C_{CE} \approx C_{CE0}$
<ul style="list-style-type: none"> • Fringing/Overlap capacitances: 	$C_{BC0}, C_{BE0}, C_{CE0}$
<ul style="list-style-type: none"> • Junction capacitance: <ul style="list-style-type: none"> – <i>Reverse biased diode</i> – <i>M=1/2-1/3, from doping</i> – <i>Φ approx 0.8V</i> – <i>M, Φ usually determined from two bias points</i> 	$C_j = \frac{C_0}{\left(1 - \frac{V_{BE}}{\Phi}\right)^M}$
<ul style="list-style-type: none"> • Diffusion capacitance <ul style="list-style-type: none"> – <i>from transit time</i> 	$C_D = g_m \tau_T \quad , \quad \tau_T = 1/2\pi f_T$

An important figure of merit: the BJT Current gain and the transit frequency f_T



The transistor is connected as a current amplifier:

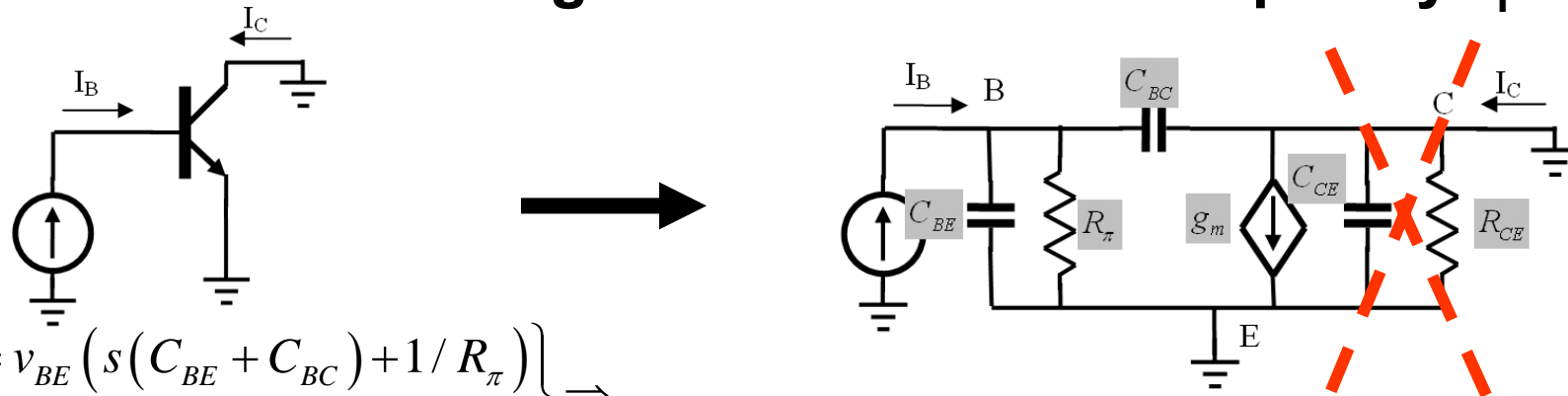
- Driven by a current source
- Driving a current meter

f_T is the frequency at which the current gain drops to unity.

f_T is usually specified by the transistor manufacturer as a **figure of merit**.

f_T is a rough estimate of the highest frequency at which the transistor can be used as an amplifier. Typically $f_T / 10$ is this highest frequency.

An important figure of merit: the BJT Current gain and the transit frequency f_T



$$\left. \begin{aligned} i_B &= v_{BE} \left(s(C_{BE} + C_{BC}) + 1/R_\pi \right) \\ i_C &= g_m v_{BE} \end{aligned} \right\} \Rightarrow$$

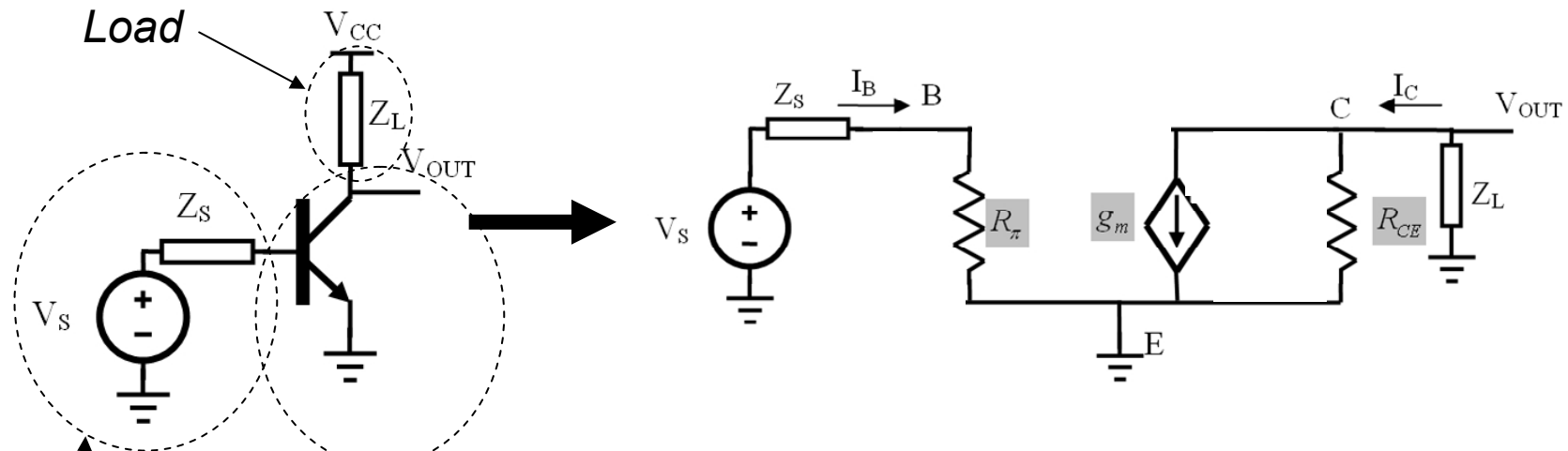
$$\frac{i_C}{i_B} = "h_{fe}" = \frac{g_m}{s(C_{BE} + C_{BC}) + 1/R_\pi} = \frac{g_m R_\pi}{1 + sR_\pi(C_{BE} + C_{BC})} = \frac{\beta_0}{1 + s\beta_0 / 2\pi f_T}$$

where $f_T = \frac{\beta_0}{2\pi R_\pi (C_{BE} + C_{BC})} = \frac{g_m}{2\pi (C_{BE} + C_{BC})} \approx \frac{g_m}{2\pi C_{BE}}$, since: $R_\pi = \frac{\beta_0}{g_m}$

the current gain is $h_{fe}(f) = \frac{\beta_0}{1 + j\beta_0 f / f_T} \Rightarrow \boxed{\lim_{f > f_T/\beta_0} |h_{fe}| = \frac{f_T}{f} \text{ with } f_T = \frac{1}{2\pi\tau_T}}$

f_T allows the easy calculation of $C_{BE} + C_{BC}$ since g_m only depends on the collector bias current! Note that $C_{BE} \gg C_{BC}$ so knowing f_T defines C_{BE}

The Common Emitter amplifier

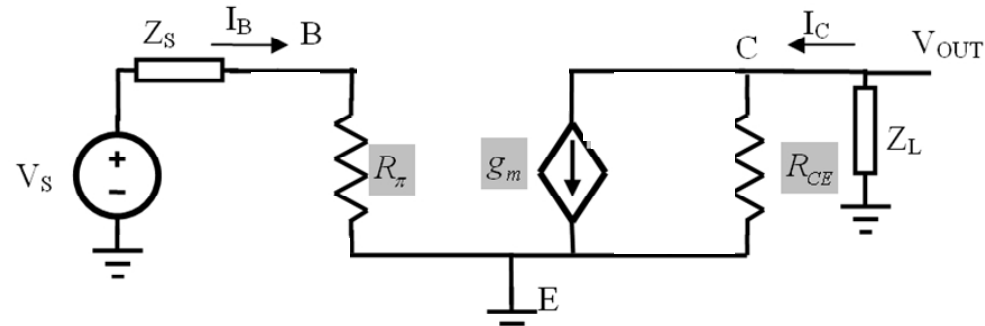


Model parameters:

$$\left\{ \begin{aligned}
 g_m &= \frac{\partial I_C}{\partial V_{BE}} = \frac{I_C}{V_T} \\
 R_\pi &= \left(\frac{\partial I_B}{\partial V_{BE}} \right)^{-1} = \left(\frac{\partial (I_C / \beta)}{\partial V_{BE}} \right)^{-1} = \beta \left(\frac{\partial I_C}{\partial V_{BE}} \right)^{-1} = \frac{\beta}{g_m} \\
 R_{CE} &= \left(\frac{\partial I_C}{\partial V_{CE}} \right)^{-1} \approx \frac{V_A}{I_C}
 \end{aligned} \right.$$

$$V_T = kT / q = 25\text{mV at } 290\text{K} = 17\text{C}$$

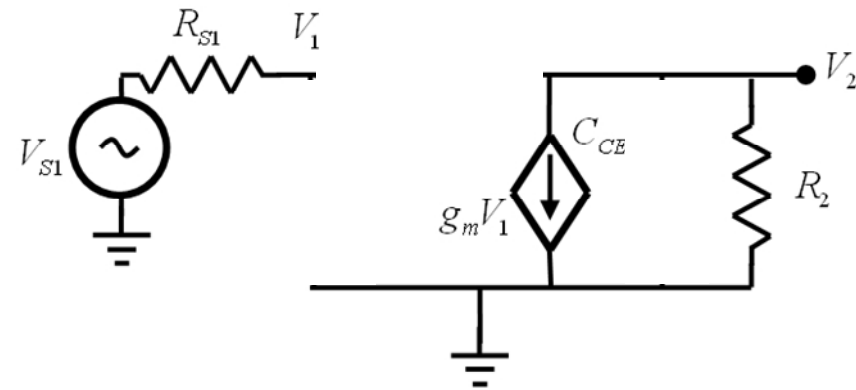
The Common Emitter amplifier (2)



Combine V_S, Z_S, R_π into a Thevenin circuit:

$$V_{S1} = V_S \frac{R_\pi}{R_\pi + Z_S} = \frac{V_S}{1 + Z_S Y_{in}}$$

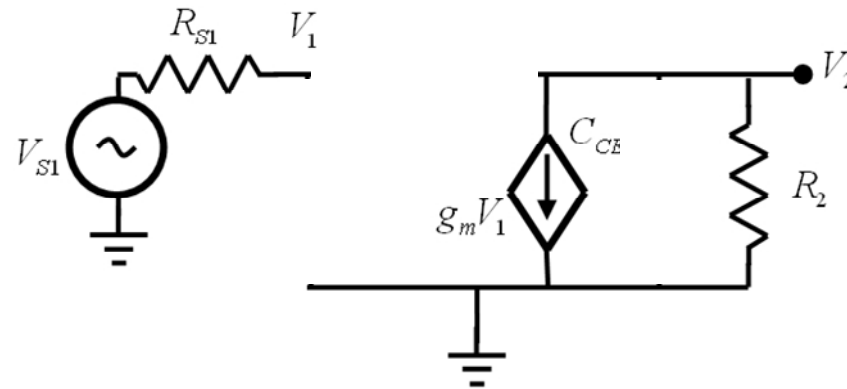
$$R_{S1} = R_\pi // Z_S = \frac{R_\pi Z_S}{R_\pi + Z_S} = \frac{Z_S}{1 + Z_S Y_{in}}$$



Consider R_{CE} and Z_L together: $R_2 = R_{CE} // Z_L = \frac{R_{CE} Z_L}{R_{CE} + Z_L} = \frac{Z_L}{1 + Z_L Y_{out}}$

and note that $Z_S = R_S + R_{BB}$

The Common Emitter amplifier (3)



Gain:

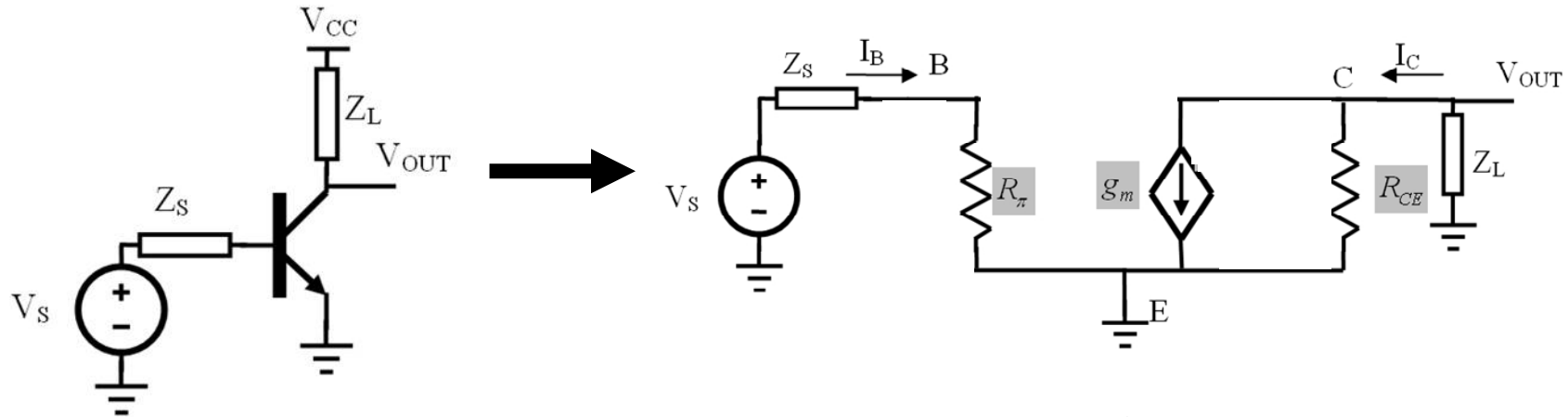
$$A_v = \frac{v_{out}}{v_s} \equiv \frac{\partial V_2}{\partial V_s} = -\frac{V_{s1}}{V_s} g_m R_2 = \frac{1}{(1 + R_s Y_{in})} \frac{-g_m Z_L}{(1 + Z_L Y_{out})}$$

If $Y_{in} = 0$ and $Y_{out} = 0$ this is the familiar:

$$A_v = \frac{v_{out}}{v_s} \equiv -g_m Z_L$$

Note that lowercase letters represent small signal variations, i.e. differentials

The Common Emitter amplifier (4): Input and output impedance



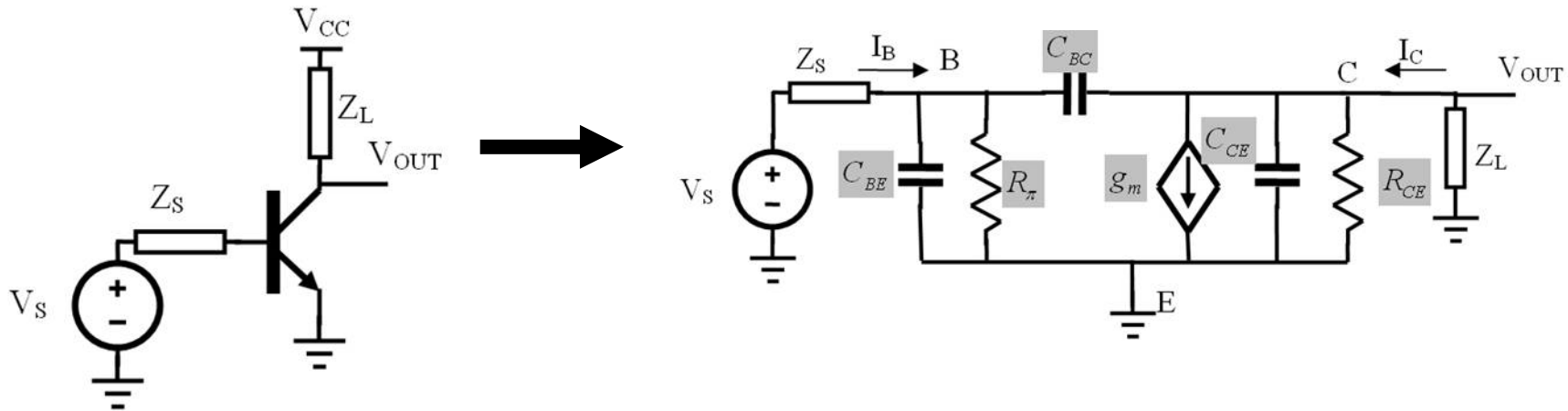
Input Impedance:

$$Z_{in} = \frac{\partial V_{BE}}{\partial I_B} = \left(\frac{\partial I_B}{\partial V_{BE}} \right)^{-1} = R_\pi$$

Output Impedance:

$$Z_{out} = \frac{\partial V_C}{\partial I_{out}} = \left(\frac{\partial I_C}{\partial V_{CE}} \right)^{-1} = R_{CE}$$

Frequency response of the Common Emitter amplifier



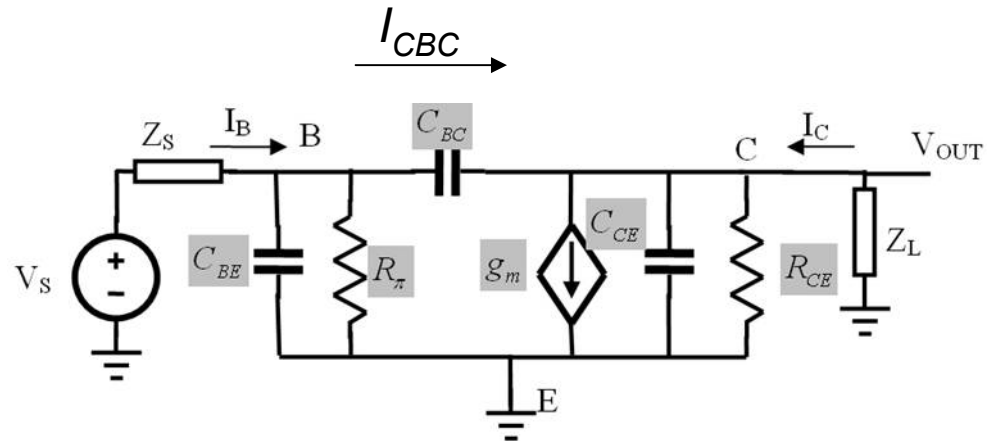
If we momentarily neglect C_{BC} this would be the same as before but with:

Input Impedance:
$$Z_{in} = \frac{1}{Y_{in}} = R_{\pi} // C_{BE} = \frac{R_{\pi}}{1 + j\omega C_{BE} R_{\pi}}$$

Output Impedance:
$$Z_{out} = \frac{1}{Y_{out}} = R_{CE} // C_{CE} = \frac{R_{CE}}{1 + j\omega C_{CE} R_{CE}}$$

We still have:
$$A_V = \frac{v_{out}}{v_s} = \frac{1}{(1 + R_s Y_{in})} \frac{-g_m Z_L}{(1 + Z_L Y_{out})} \quad \text{and} \quad \frac{v_{out}}{v_B} = G = \frac{-g_m Z_L}{(1 + Z_L Y_{out})}$$

Dealing with C_{BC}



In the absence of C_{BC} we know that: $\frac{v_{out}}{v_B} = G = \frac{-g_m Z_L}{(1 + Z_L Y_{out})}$

the current through C_{BC} is:

$$I_{CBC} = j\omega C_{BC} (V_B - V_C) = j\omega C_{BC} (1 - G) V_B = j\omega C_{BC} (1/G - 1) V_C$$

We can account for C_{BC} by introducing an additional input “Miller” admittance:

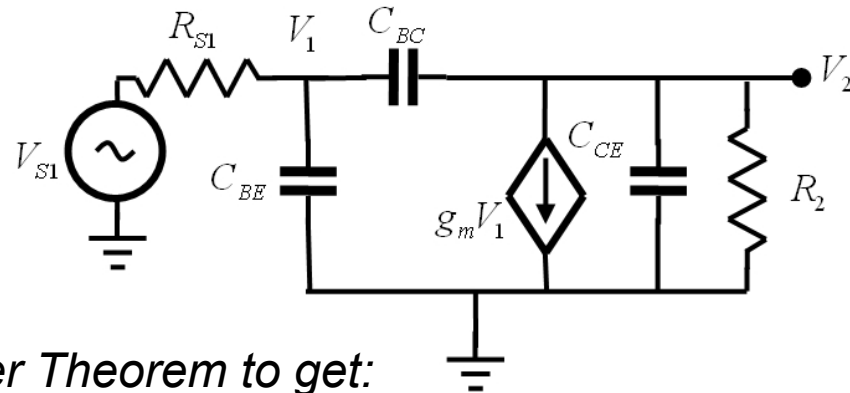
$$Y_{M,in} = j\omega C_{BC} (1 - G)$$

and an additional output “Miller” admittance:

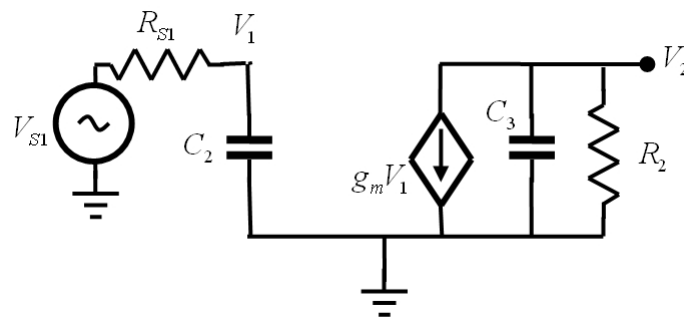
$$Y_{M,out} = j\omega C_{BC} (1 - 1/G)$$

Frequency response of the CE amplifier

V_{S1} and R_{S1} is the Thevenin equivalent of R_S , V_S and R_π



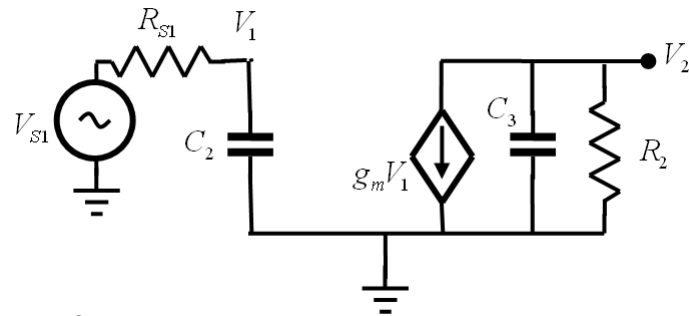
Use the Miller Theorem to get:



with $C_2 = C_{BE} + (1 + G)C_{BC} \approx (1 + G)C_{BC}$ and $C_3 = C_{CE} + (1 + 1/G)C_{BC} \approx C_{BC}$

The low frequency gain from V_1 to V_2 is: $G = -g_m R_2 \approx -g_m Z_L$

Frequency response of the CE amplifier (2)



the gain of this amplifier is:

$$\frac{dV_2}{dV_{S1}} = \frac{-g_m}{1 + sR_{S1}C_2} \frac{R_2}{1 + sR_2C_3} = \frac{G}{\underbrace{(1 + sR_{S1}C_2)}_{\text{input pole}} \underbrace{(1 + sR_2C_3)}_{\text{output pole}}}$$

If the gain is large the input pole can be at a much lower frequency than the output pole:

$$\tau_{IN} = R_{S1}C_2 = (R_\pi // R_S)(C_{BE} + (1+G)C_{BC}) \approx R_S(1+G)C_{BC} \quad (\text{since } R_S \ll R_\pi)$$

$$\tau_{OUT} = R_2C_3 = (C_{BC}(1+1/G) + C_{CE})(R_{CE} // Z_L) \approx C_{BC}Z_L \quad (\text{since } C_{BC} \gg C_{CE})$$

if $\tau_{in} > \tau_{out}$ The gain-bandwidth product of this amplifier is:

$$\left. \begin{aligned} |G| &= g_m R_2 \approx g_m Z_L \\ BW &= \frac{1}{2\pi R_{S1}C_2} \approx \frac{1}{2\pi R_S C_2} \end{aligned} \right\} \Rightarrow GBW = \frac{g_m R_2}{2\pi R_S C_2} \cdot \lim_{|G| \rightarrow \infty} GBW \approx \frac{G}{2\pi R_S C_{BC}(1+G)} \approx \frac{1}{2\pi R_S C_{BC}}$$

This is independent of G!

The tiny C_{BC} and the **source impedance** can define the gain-bandwidth product of the CE amp.

This suggests that feedback is at work!

Note that the input pole will be at a lower frequency than the output pole if

$$|G|R_{S1} > R_2 \Rightarrow g_m R_{S1} R_2 > R_2 \Rightarrow g_m R_S = \beta R_S / R_\pi > 1$$

Importance of the gain-bandwidth product

the gain of the CE amplifier is:

$$\frac{dV_2}{dV_{s1}} = \frac{-g_m}{1 + sR_{s1}C_2} \frac{R_2}{1 + sR_2C_3} = \frac{G}{\underbrace{(1 + sR_{s1}C_2)}_{\text{input pole}} \underbrace{(1 + sR_2C_3)}_{\text{output pole}}}$$

One of the two pole frequencies is usually much lower than the other.

*We then say that we have a **dominant pole amplifier** whose frequency response is:*

$$A_v(s) = \frac{G_0}{1 + s\tau} = \frac{G_0}{1 + j\omega\tau} \quad \text{for a DC gain } G_0 \text{ and a bandwidth } \omega_B = 1/\tau$$

We have already seen that in the limit of large G_0 **the Gain-Bandwidth product:**

$$G_B = G_0\omega_B = G_0 / \tau$$

is independent of the amplifier gain.

We will see later in the course that this is a very general result which **applies to all one pole feedback amplifiers** regardless of how they are constructed.